**Stress-Strain Analysis of Cubic Crystals**

Maninder Kaur,

PG Department of Physics,

DAV College,

Katra Shar Singh, Amritsar,

Punjab, India.

mannu\_711@yahoo.co.in

**I. Introduction**

Condensed matter physicsis the branch of physics that is concerned with the microscopic and macroscopic physical properties of the condensed phases of matter.We study condensed matter physics because it explains the properties of solids and fluids in our day-to-day life. Most of the objects around us are in the condensed phase and all of these have different properties. Metals are lustrous, glass is transparent and fluids wet the surface in contact with them. The reason behind the different properties of these materials is the underlying arrangements of their constituent atoms, which is well explained by condensed matter physics.

Finally, it is a very worthwhile branch to cater to our need for new materials. Over the last few years, with the help of condensed matter physics, we have engineered new materials to revolutionize our world completely. The most prominent example is the progress in semiconductor technology, which is the backbone of computers, iPhones, and the whole e-world.

1. **Motivation behind this chapter**

We need to know Bulk Modulus, Poisson’s Ratio, Young’s Modulus, and Rigidity Modulus for engineering work. These are expressed in terms of the elastic compliance constant and elastic stiffness constant. The relation between elastic compliance constant and elastic stiffness constant has been derived only partially in all the leading textbooks like R.L. SINGHAL, A.J. DEKKER & CHARLES KITTEL, etc. This chapter aims to fully derive the relation between the elastic compliance constant and elastic stiffness constant to facilitate the understanding of this topic of the Post Graduate Condensed Matter Physics 1.

1. **Stress-strain analysis**

 Stress-Strain Analysis aims to determine the stresses and strains in materials and structures when they are under the effect of some force. As a reaction to the applied force, a restoring force emerges in the material. Stress is defined as the restoring force per unit area, produced in the body against deformation.

Thus, Stress is the ratio of force per unit area i.e., S=R/A, where S → Stress, R → Internal resisting force, and A → cross-sectional area.

 Moreover, the applied force has the tendency to change the dimensions as well as the shape of the body. This change is measured in terms of a physical quantity known as Strain. Strain is defined as the ratio of change in dimensions to the original dimension of the body when it is subjected to some external force.

1. **General principle**

In condensed matter physics, the target of stress analysis is to evaluate the stress-strain values for solid objects. It is not applicable to fluids as it is. The study of stresses in fluids is concerned with fluid mechanics. Stress analysis considers the bulk portion of materials and the characteristics of continuum mechanics. The bulk is further imagined to be made of small cubes which contain a large number of atoms, and its properties are averages of the properties of those atoms. In stress analysis, we ignore the cause of forces or the nature of the materials undertaken. We also assume the stresses and the resulting strain in the materials are related through known mathematical equations.

**II. Elastic Constants**

When a material, considered as an elastic body is subjected to a force, stress is built up in the material. This stress generates a strain on the material. The ratio of the applied stresses to the strains produced are always constant. Further, there are two types of stress i.e., Normal Stress (tensile or compressive) and Shear Stress. Similarly, strain is of four types; longitudinal strain, lateral strain, volumetric strain and shear strain. The ratio of specific stress and strain is represented in terms of quantities known as Elastic constants. There are four types of elastic constants: -

1. Young’s Modulus
2. Bulk Modulus
3. Rigidity Modulus
4. Poisson’s Ratio
5. **Young’s Modulus**

According to Hook’s Law, when a body is subjected to tensile stress ( where the applied force is perpendicular to the cross-section area of the body and results in an increase in the length,), the strain produced is directly proportional to the stress within the elastic limits of that body. The ratio of applied stress to the strain is constant and is known as Young’s modulus.



1. **Bulk Modulus**

When equal and identical, mutually perpendicular normal stresses are applied to a body, within its elastic limits, volumetric strain is produced. Bulk modulus is defined as the ratio of direct stress to the corresponding volumetric strain.



1. **Rigidity Modulus**

When a body is subjected to shear stress (where the force applied is along the area of cross-section), it results in shear strain i.e., the shape of the body changes. The ratio of shear stress to the resulting shear strain is called rigidity modulus.



1. **Poisson’s Ratio**

When a simple tensile stress is applied to a body within its elastic limits, then there is a change in the measurements of the body in the direction of the load as well as in the perpendicular direction. When these changed measurements are compared with their original ones, longitudinal strain, and lateral strain are obtained. The ratio of the Lateral strain to the Longitudinal strain is called Poisson’ ratio.



 **Iii. Components Of Stress**

Let us Imagine a small cube of sides dx, dy, and dz removed from the solid, on which stress is applied. On each face of the cube, the forces can be resolved into three mutually perpendicular components, one is normal to face which results in normal stress and two are lying in the plane of the face that form shear stress.



**Figure 1: Components of force**

Here fy results in normal stress, and both fx and fz result in shear stress. However, when the cube is in dynamic equilibrium, forces on opposite faces must be equal in magnitude and opposite in sign.



**Figure 2:** **Dynamic Equilibrium**

Thus, to describe the stress condition of the cube Nine components are required. Three stresses are normal to cube faces and Six stresses act across the cube faces.



**Figure 3: Components of stress**

Stress components in matrix form are,



**Figure 4: Stress Matrix**

Where the Uppercase alphabets indicate the direction of force and the subscripts denote the direction of normal to the plane on which force is acting. The number of independent stress components are further reduced to six when it is a cube in static equilibrium, i.e., the cube does not rotate.

* Under this condition

 Xy = Yx ; Zy  = Yz ; Zx = Xz

* The six independent stress components then are, Xx, Yy, Zz, Xy, Yz, Zx



**Figure 5: Static Equilibrium**

**IV. Components of strain**

 Consider an unstrained solid with orthogonal unit vectors *i, j, and k* as shown in Figure 6. For an orthogonal system, Δa, Δb, Δc, Δα, Δβ, Δϒ correctly define the six components of elastic strain. Where α,β, ϒ is the angle between the unit cell axes a,b,c.



**Figure 6: Orthogonal System**

 However, for the non-orthogonal axis, this leads to mathematical complications. Hence, a general situation strain is specified with the use of six elements exx, eyy, ezz, exy, eyz, and *ezx* as described afterward.

 Suppose a small uniform deformation i.e., deformation in which each primitive cell of crystal is deformed in the same way, of the solid that results in distorted orientation and length of axes as shown in figure.7,



**Figure 7: Non-Orthogonal System**

The new axis *i’, j’, k’* is expressed in terms of the old axis as given below

i’ = (1+ϵxx)i + ϵxyj + ϵxzk …….1(a)

j’ = ϵyxi + (1+ϵyy)j + ϵyzk ……..1(b)

k’ = ϵzxi + ϵzyj + (1+ϵzz)k ………1(c)

where the coefficients ϵxx, ϵxy, ϵxz, etc. identify the deformation; and are dimensionless quantities with values that are very less than unity, i.e. the strain is meagre. Also, the fractional changes in the length *of* axis *i,* to the first order, is *ϵxx,* and that of *j* and *k* axes are *ϵyy* and *ϵzz.* Thus *ϵxx, ϵyy, ϵzz* represents linear strain components which are defined as

 exx = ϵxx, eyy = ϵyy, ezz = ϵzz  *……(2)*

Similarly*,* i’. j’ ≅ ϵyx + ϵxy, which gives a measure of change in orientation between *i’& j’* due to stress. Thus, *exy* = change in angle between axes *i’& j’.*

Thus

 exy = i’. j’ ≅ ϵyx + ϵxy

 eyz = j’. k’ ≅ ϵyz + ϵzy

 ezx = k’.i’ ≅ ϵzx + ϵxz  *..….(3)*

Now merely rotating the axes does not change the angle between them. So using equation *(3),* a pure rotation represented by ϵyx = -ϵxy, ϵzy = -ϵyz & ϵzx = -ϵxz are excluded. Further taking ϵyx = ϵxy, ϵzy = ϵyz & ϵzx = ϵxz *.*values in equation *(3)* , we get

exy = 2ϵxy; eyz = 2ϵyz; ezx = 2ϵzx  *……(4)*

 **V. RELATION BETWEEN STRESS ANS STRAIN COMPONENTS**

According to Hook’s Law, when tensile stress is applied to a body, the strain produced is directly proportional to the stress under the elastic point of that body. Hence, the strain components can be expressed as linear functions of stress components :



 *……(5)*

Conversely, the stress components are expressed as linear functions of the strain components:



 *……(6)*

 The coefficients **S11, S12,** and extra are termed as **ELASTIC COMPLIANCE CONSTANT,** and the coefficients **C11, C12** and extra are termed as **ELASTIC STIFFNESS CONSTANT.** The Elastic Compliance Constants have the dimensions of [strain/stress]= [1/(force/area)]= [area/force] or [area x length/force x length]= [volume/energy]. TheElastic Stiffness Constants have the dimensions of [stress/strain] =[force/area] or [energy/volume].

**VI. DERIVATION OF STRESS-STRAIN COMPONENTS FOR CUBIC CRYSTAL**

We will prove that cubic crystals have only three independent stiffness constants. We begin with the declaration that the elastic energy density of a cube crystal is given as

U = 1/2C11(e2xx + e2yy + e2zz) + 1/2C44(e2yz + e2zx + e2xy) + C12(eyyezz + ezzexx + exxeyy)

 ……(7)

 and it does not include any other quadratic terms like

 (exxexy + ……..); (eyzezx + ……..); (exxeyz + ………); ………. ……(8)

 This result is a consequence of the minimum symmetry requirement of a cubic crystal, which is based on th four three-fold rotation axes passing through the body’s diagonal directions. One such rotation axis is shown in Figure 7. If we rotate by an angle of 120֯ about this body diagonal the x-axis changes to y, the y-axis changes to z, and z -axis changes to x. Replacing x by y, y by z, z by x in equation (7) gives

 U = 1/2C11(e2yy + e2zz + e2xx) + 1/2C44(e2zx + e2xy + e2yz) + C12(ezzexx + exxeyy + eyyezz)

Hence, equation (7) does not vary under the operation considered.

**

 **Figure 7: Rotation by 2π/3 about body diagonal**

Hence, equation (7) ) does not vary under the operation considered.Similar rotations about the other axis of rotations change:

[ -x-axis changes to z, z-axis changes to -y and -y -axis changes to -x];

 [ x-axis changes to z, z -axis changes to -y and -y -axis changes to x];

 [ -x -axis changes to y, y-axis changes to z and z -axis changes to -x].

 These rotations will again leave equation (7) invariant. However, if the odd terms as in equation (8) are included in the energy density term, it may result in a change of sign after the rotation operations. E.g., *exy = -ex(-y).* Hence the terms in equation (8) are not invariant under the rotation operations, and hence should not be considered in the energy density expression. This proves that our assertion is true.

 Now, differentiating equation (7) w.r.t*. exx,* we get

$$\frac{∂U}{∂e\_{xx}}=C\_{11 }e\_{xx}+ C\_{12}(e\_{yy}+ e\_{zz})$$

 But,

 $\frac{∂U}{∂e\_{xx}}= X\_{x}$

 Hence $, $

$$X\_{x}=C\_{11 }e\_{xx}+ C\_{12}(e\_{yy}+ e\_{zz})$$

Comparing it with,

$$X\_{x}=C\_{11 }e\_{xx}+ C\_{12}e\_{yy}+C\_{13}e\_{zz}+ C\_{14}e\_{yz}+ C\_{15}e\_{zx}+ C\_{16}e\_{xy}$$

 we get,

*C12 = C13 and* *C14 = C15 =C16 = 0 …….(9)*

Similarly Comparing

$X\_{y}=$ $\frac{∂U}{∂e\_{xy}}=C\_{44 }e\_{xy}$

With,

$$X\_{y}=C\_{61 }e\_{xx}+ C\_{62}e\_{yy}+C\_{63}e\_{zz}+ C\_{64}e\_{yz}+ C\_{65}e\_{zx}+ C\_{66}e\_{xy}$$

we get,

*C61 = C62 = C63 = C64 = C65 = 0 and C66 = C44 ……(10)*

Proceeding with the similar calculation we find that array depicted by equation (6) takes the form of a simpler matrix for a cubic crystal as given below,

**……..(11)

This proves our declaration that for a cubic crystalC11, C12, and C44are the only three independent elastic stiffness constants*.* Here constants C11 relate the compression stress and strain along the *X,* *Y,* or *Z* axis, while *C44* relates the shear stress applied to a certain direction to the strain in the same direction i.e*., Yz = C44eyz, Zx = C44ezx* and so on. The constant C12 relates the compression stress in one direction to the strain in other perpendicular directions i.e., *eyy* with *Xx, ezz* with *Xx, exx* with *Yy & Zz, ezz* with *Xx & Yy* as seen from matrix (11)*.*

 Similarly, the inverse matrix matrix takes the form as;

 ……..(12)

Now from matrix (11)

$$X\_{x}=C\_{11 }e\_{xx}+ C\_{12}e\_{yy}+C\_{12}e\_{zz}$$

$$Y\_{y}=C\_{12 }e\_{xx}+ C\_{11}e\_{yy}+C\_{12}e\_{zz}$$

$$Z\_{z}=C\_{12 }e\_{xx}+ C\_{12}e\_{yy}+C\_{11}e\_{zz}$$

*.…..(13)*

Also from matrix (12)

$$e\_{xx}= S\_{11}X\_{x}+ S\_{12}Y\_{y}+ S\_{12}Z\_{z}$$

 *.…..(14)*

From equation (13) and (14)

$e\_{xx}= S\_{11}\left(C\_{11 }e\_{xx}+ C\_{12}e\_{yy}+C\_{12}e\_{zz}\right)+ S\_{12}\left(C\_{12 }e\_{xx}+ C\_{11}e\_{yy}+ C\_{12}e\_{zz}\right)+ S\_{12}($ $C\_{12 }e\_{xx}+ C\_{12}e\_{yy}+C\_{11}e\_{zz})$ *……(15)*

$e\_{xx}= \left(S\_{11 }C\_{11}+ S\_{12}C\_{12}+S\_{12}C\_{12}\right)e\_{xx}+\left( S\_{11}C\_{12}+ S\_{12}C\_{11}+S\_{12}C\_{12}\right)e\_{yy}+ ($ $S\_{11 }C\_{12}+ S\_{12}C\_{12}+S\_{12}C\_{11})e\_{zz}$  *……(16)*

Equating coefficients on L.H.S. & R.H.S.

 $\left(S\_{11 }C\_{11}+ S\_{12}C\_{12}+S\_{12}C\_{12}\right)$ = 1 *.…..(17)*

 $\left( S\_{11}C\_{12}+ S\_{12}C\_{11}+ S\_{12}C\_{12}\right)$ = 0 *……(18)*

Rearranging equation *(18)*

 $S\_{11}= \frac{-S\_{12}(C\_{11}+ C\_{12}) }{C\_{12}}$  *.…..(19)*

Putting equation *(19)* in *(17)*

$\frac{-S\_{12}(C\_{11}+ C\_{12}) }{C\_{12}}+ 2S\_{12}C\_{12}=1$ *…...(20)*

$$S\_{12}\left[2C\_{12}- ^{C\_{11}}/\_{C\_{12}}\left(C\_{11}+C\_{12}\right)\right]=1$$

$$-S\_{12}\left[C\_{11}^{2}+C\_{11}C\_{12}-2C\_{12}^{2}\right]=C\_{12}$$

$$-S\_{12}\left[C\_{11}^{2}+2C\_{11}C\_{12}-C\_{11}C\_{12}-2C\_{12}^{2}\right]=C\_{12}$$

$$-S\_{12}\left[ C\_{11}\left(C\_{11}+2C\_{12}\right)-C\_{12}(C\_{11}+2C\_{12})\right]=C\_{12} $$

 $S\_{12}=\frac{-C\_{12}}{(C\_{11}-C\_{12})(C\_{11}+2C\_{12})}$ *.…..(21)*

Putting equation *(21)* in *(19)*

$$S\_{11}=\frac{(-1)(-C\_{12})}{(C\_{11}-C\_{12})(C\_{11}+2C\_{12})} \left(\frac{C\_{11}+C\_{12}}{C\_{12}}\right)$$

$$S\_{11}=\frac{C\_{11}+C\_{12}}{(C\_{11}-C\_{12})(C\_{11}+2C\_{12})}$$

  *…… (22)*

Similarly from matrix *(11)*

$Y\_{z}=e\_{yz }C\_{44};$ $Z\_{x}=e\_{zx }C\_{44};$ $X\_{y}=e\_{xy }C\_{44}$.

and from matrix *(12)*

$e\_{yz}=S\_{44 }Y\_{z};$ $e\_{zx}=S\_{44 }Z\_{x};$ $e\_{xy}=S\_{44 }X\_{y}.$

Which leads to,

$$Y\_{z}=C\_{44 }S\_{44} Y\_{z}$$

$$S\_{44}=\frac{1}{C\_{44 }}$$

  *……(23)*

From matrix *(12),* it is clear that tensile stresses *Xx, Yy, Zz* does not produce any shear strain i.e. *eyz = ezx =exy = 0* for *Xx, Yy, Zz.* Also pure shear stresses *Yz, Zx, Xy* produce the pure shear strain only i.e. $e\_{yz}=S\_{44 }Y\_{z};$ $e\_{zx}=S\_{44 }Z\_{x};$ $e\_{xy}=S\_{44 }X\_{y}$. Thus it is concluded that terms S11  gives the amount of tensile strain in x-direction generated due to tensile stress of unit magnitude along the x-axis. And term S12 gives the extent of y-axis or z-axis strain tha comes out from the x-axis tensile stress.

 The two-dimensional view is like,



 

Thus, S12 is equivalent to the lateral strain which acts perpendicular to the direction of applied stress.

1. **KNOWING THE PROPERTIES OF CUBIC CRYSTALS IN TERMS OF STRESS-STRAIN COMPONENTS**

 We know that Young’s Modulus or modulus of elasticity is expressed as the fraction of tensile or compressive stress and the resultant strain i.e.

*Young’s Modulus(E) = Tensile Stress/Tensile Strain*

 Hence, it is clear that

…… *(24)*

Similarly, Poisson’s ratio is expressed as the fraction of lateral strain and longitudinal strain i.e.

 *Poissons’s ratio (µ) = lateral strain/longitudinal strain*

Hence, it is clear that.

…..*. (25)*

Also, the shear modulus or modulus or modulus of rigidity G is expressed as the fraction of shear stress and shear strain i.e.

 *Rigidity Modulus (G) = Shear Stress/Shear Strain*

 Hence, it is clear that

….*..(26)*

 Further the bulk modulus is expressed in terms of the the change in volume of a body as a result of unit compressive or tensile stress.

 For this condition a uniform dilation of crystal:

$$e\_{xx}=e\_{yy}=e\_{zz}=\frac{1}{3}δ$$

 Hence,

$$U=\frac{1}{2}C\_{11}\left[\frac{1}{9}δ^{2}+\frac{1}{9}δ^{2}+\frac{1}{9}δ^{2}\right]+C\_{12}\left[\frac{1}{9}δ^{2}+\frac{1}{9}δ^{2}+\frac{1}{9}δ^{2}\right]$$

*U* = $\frac{1}{2}C\_{11}\left[\frac{3}{9}δ^{2}\right]+C\_{12}\left[\frac{3}{9}δ^{2}\right]$

*U* = $\frac{1}{6}C\_{11}δ^{2}+\frac{1}{3}C\_{12}δ^{2}$

*U* = $\frac{1}{6}(C\_{11}+2C\_{12})δ^{2}$

Now the Bulk Modulus is produced by by the equation.

$$U= \frac{1}{2}Bδ^{2}$$

Hence,

$B=\frac{1}{3}(C\_{11}+C\_{12})$ .*…..(27)*

Further, the compressibility K is the defined as

$K=\frac{1}{B}=\frac{3}{C\_{11}+2C\_{12}}$ *……(28)*

If *C12=*0, from equations *(24), (25), (27) & (28),* We get

E=C11, µ=0, B=1/3 C11 & K=3/C11

Which happens when there is no transverse contraction even when there exists a longitudinal expansion. This approximation is useful when one is handling a one-dimensional lattice.

 Further, if the crystal is an isotropic,

 C12 + 2C44 = C11 [ or S44 = 2(S11 – S12)]

 i.e. there are only two independent moduli.

* ***APPLICATIONS: -***
* Stress-Strain Analysis is important when we design the constructions such as tunnel, bridges and dams etc.
* It is primarily used for civil mechanics and aerospace engineers.
* It is used for theupkeep constructions and to diagnose the reasons of structural failures.
* ***REFERENCES: -***
1. Introduction Solid State Physics by CHARLES KITTEL
2. Elements of Solid-State Physics by J.P. SRIVASTAVA
3. http://en.m.wikipedia.org/wiki/Stress%E2%80%93strain\_analysis

1. Introduction to Solid State Physics by ARUN KUMAR
2. Solid State Physics by A.J. DEKKER
3. Solid State Physics by R.L. SINGHAL
4. Solid State Physics by M A WAHAB