# A unique potential to study scattering and fusion phenomena in heavy ion collisions around Coulomb barrier

### Bidhubhusan Sahu1*,*∗, Kamala Kanta Jena2, Santosh Kumar Agarwalla3

Department of Physics, School of Applied Sciences, KIIT Deemed to be University, Bhubeneswar-751024, India

2P. G. Department of Physics, Bhadrak Autonomous College, Bhadrak-756100, India

3P. G. Department of Physics, Fakir Mohan University, Balasore-756019, India

∗Emails: [bbsahufpy@kiit.ac.in,](mailto:bbsahufpy@kiit.ac.in) [kkjena1@gmail.com,](mailto:kkjena1@gmail.com) [san1612@rediffmail.com](mailto:san1612@rediffmail.com)

**Abstract:** To examine angular variation in cross-sections of elastic scattering, a systematic recursive formula is developed involving partial-wave scattering matrix for total effective complex potential of nucleus-nucleus collisions. We express cross-sections for the absorption from arbitrary and small intervals in an additional manner. This results in an assessment of the effective potential’s absorption contributions in the interior region, which explains fusion cross section (*σfus*) data for different incident energies. The interaction potential taken into account in this study is energy-independent and enables states of different partial-wave trajectories owing to its weakly absorbing characteristic. Therefore, it becomes clear that the appearance of these resonance states is the physical cause of the observable oscillatory structure in the modification of the quantity D(E*c.m.*), the second derivative of the product “E*c.m.σfus*” with respect to E*c.m.*. In this chapter, we discuss simultaneous and extremely effective descriptions of the crosss-ections for fusion, elastic scattering, and the outcomes of D(E*c.m.*) in various kinds of heavy-ion collisions.

# Introduction

# There have been numerous experiments on the nucleus-nucleus collision process, and extensive data on the angular distribution of the elastic scatterings ) for various energies of incidence and the fusion cross-sections at firmly energy intervals are now available [[1]](#_bookmark16) -[[5].](#_bookmark20)Eight such systems come to mind in this regard: 12C+208Pb in Refs. [[1,](#_bookmark16) [3],](#_bookmark18)16O+208Pb in Refs. [[2,](#_bookmark17) [4],](#_bookmark19)19F+208Pb in Refs.[[6,](#_bookmark21)[7],](#_bookmark22)16O+144Sm in Refs.[[8–10],](#_bookmark25)and 16O+62Ni in Refs.[[11,](#_bookmark26)[12]](#_bookmark27)all of which include a wealth of information. The shape resonances, though present, are not observed experimentally in heavy ion systems. But the case is different for the systems with light ions, where the resonances are generated by the effective potential [[13].](#_bookmark28) Further research is needed to ascertain the potential impact and expression of resonances in any other visible forms throughout collision. One employs the phenomenological potential, which is often complicated, to analyse these heavy ion collision data. The observed values of cross section in the elastic scattering at different incident energies are reproduced to determine all the parameters defining the potential. One has to explain fusion cross-section (*σfus*) data obtained from fusion process and the phenomena of resonance taking place in the elastic scattering process by using the same interaction potential. It is thus challenging to device a special potential that can address both of these phenomena at the same time. It is because scattering process is surface-sensitive, whereas, fusing process volume-sensitive. The scattering is associated with potential’s nature on surface region, but fusion is an internal activity. Once we obtain data from elastic scattering for theoretical analysis, the results of total reaction cross section (*σr*) can be found easily. Total cross-section is the sum of cross sections for various reaction channels, in which the fusion channel predominates in low-energy collisions. It is difficult to separate the part of reaction cross section from the total value that is an exact amount of the measured data of σfus.

# The cross-section *σr* is taken into account in the scope of the optical potential model analysis of scattering by the expected value of the imaginary component of the potential, calculated using the distorted waves from the full potential in the elastic channel. This is just the sum of the cross-sections produced by absorption in all of the potential's different regions where the imaginary portion is present. We can use the same wave function as used to explain the elastic scattering data to obtained the absorption cross-section in ith infinitesimally small ith radial interval so that the total absorption cross-section can be written as , where n represents the totality of intervals of the potential with potential range . The absorption cross-section can be explicitly obtained at various intervals of the potential by using above expression. The fusion being an interior phenomena is expected to occur in the interior to the radial position of electrostatic Coulomb barrier, and the absorption in the region 0<r<RB must account for the data. The exact radius *Rfus* up-to which absorption cross section is calculated for explaining is said to be fusion radius. Udagawa et al. [[14,](#_bookmark29)[15] in their direct reaction model (DRM) use such](#_bookmark30) concept of fusion cross-section. But they have considered fusion radius greater than RB in many heavy ion collisions analyzed, which allows fusion activity to start before reaching the Coulomb barrier RB. But fusion is normally assumed to happen after full penetration of barrier. Thus the concept will move opposite to the common notion [[16,](#_bookmark31)[17]](#_bookmark32). Therefore, this was severely criticized [[18,](#_bookmark33)[19].](#_bookmark34)

# We overcome the above anomaly in our present discussion and consider the systems 12C+208Pb, 16O+208Pb,19F+208Pb, 16O+144Sm, and 16O+62Ni in order to explain the experimental results of *σexpt* by keeping *Rfus* smaller with respect to *RB*. Without employing any numerical integration method including Runge-Kutta, we use an alternative approach for solving Schro¨dinger equation for given nucleus-nucleus potentials. Our approach is suitable for investigation of interval-wise absorption in the reaction process. The potential is simulated in our computation using arbitrarily tiny rectangular pieces. We establish an analytical formula for the scattering S-matrix in order to explain data, by employing accurate wave functions and their analytical continuation between neighbouring parts. The absorption level in each tiny portion or width of the potential is determined by using the same wave functions. As a result of which the overall reaction cross-section *σr* becomes the sum of contributions due to absorption across the whole range of potential. In contrast, we consider the total contribution for absorption across a constrained region 0 *< r < Rfus* inside *RB* of the Coulomb barrier in order to explain why the fusion cross section *σfus* magnitude remains below *σr*. Interestingly, the outcomes obtained by using our approach well match results obtained from the method of Runge-Kutta.

In heavy nuclei collisions near the Coulomb barrier, measurements with excellent accuracy reveal data of *σfus* at relatively close energy intervals. For such a pair of heavy nuclei, the variation in the outcomes of *σfus*with the bombarding energy *Ec.m.* in c.m. systemshows smooth fluctuation without any distinct characteristic, but it oscillates when light pairs of nuclei collide. It the product *Ec.m.σfus* is doubly differentiated w.r.t. *Ec.m* using some point difference formula, then function shows unusual oscillatory structure as energy *Ec.m.*varies [[5].](#_bookmark20) This is referred to as barrier distribution. The experimental results of *D*(*Ec.m*) for the above cases are successfully explained in addressing peaks. The values of *σfus* obtained theoretically by using our method are presented in *D*(*Ec.m.*) form. In this study, we identify the following crucial traits in the potential that we used; in addition to being very deep and having little diffuseness, the real part is also highly strong compared to the imaginary part. Shape resonance states (experimentally not seen) [[13]](#_bookmark28) might endure throughout the collision process as a result of the formation of standing waves in the nuclear well because of the potential’s less absorptive character. The oscillating structure in the results of *D*(*Ec.m.*) as a function of *Ec.m.*is consequently attributed to these resonances.

It should be noted that coupled-channels (CC) formulation is the natural language for investigating fusion processes at energies around the Coulomb barrier. For this, a number of computer programs, including CCFUS [[20,](#_bookmark35)[21]](#_bookmark36)and CCFULL[[22],](#_bookmark37)have been created. Since there are many channels present in the heavy ion collision process due to its complexity, solving coupled equations that take all of these channels into account is both difficult and time-consuming. These are nearly schematic formulations. This method includes key approximations in order to simplify the calculation process. In the majority of pairings of nuclei, the aspects of both *σfus* and *D*(*Ec.m.*) cannot be explained satisfactorily even with exhaustive CC calculation. This unsatisfactory circumstance is still present in the most current CC projections [[23]](#_bookmark38) based on M3Y plus repulsion potential used for analyzing data of the 16O+208Pb system.

According to a most recent calculation [[24],](#_bookmark39) even observed value of *σfus* from the sub-barrier region through above barrier region cannot be concurrently estimated by CC calculations with identical the nuclear potential of Woods-Saxon. The CC calculation in microscopic level is essentially one dimensional model of crossing the barrier [[25,](#_bookmark40)[26]](#_bookmark41) that includes lots of barriers with different heights that are produced as a result of coupling between the relative motion and the internal degrees of freedom of the colliding nuclei, such as static deformation,collective vibration [[27],](#_bookmark42) inelastic excitation, and nucleon transfer [[23].](#_bookmark38) For imposing boundary condition for the barrier crossing model, the coupled-channel calculation for cross section of fusion scarcely takes the impact of any mechanism of the interaction potential in the interior pocket into account because it is believed to be exceedingly absorptive.

The potential considered for nucleus-nucleus interaction in pocket area is a key component of the current formulation and is used to explain the experimental results of *σfus* as well as the related function *D*(*Ec.m.*) generated from measured *σfus*. Due to the pocket’s less absorptive character, the resonances that it produces successfully depict the oscillatory structure of *D*(*Ec.m.*). Further more, although it is not stated directly in our model, the impact of coupling is implied. It is anticipated that the entrance channel potential barrier would vary dramatically when a non-elastic channel is coupled with it [[28],](#_bookmark43) especially in region *r < RB*, where the effective potential will suddenly drop [[29,](#_bookmark44)[30].](#_bookmark45) By choosing small diffuseness parameter in the Woods-Saxon potential, this coupling effect may be easily included in the formulation. The current formulation includes the influence of coupling of channel in a phenomenological fashion by using such a modest diffuseness value in the analysis of scattering cross-sections and fusion cross-sections simultaneously. We discuss in section 1.2 the formulation regarding analytical expression for the S-matrix and region-wise absorption. We apply our theory in section 1.3 to explain bserved data of , *σfus*(*Ec.m.*) and *D*(*Ec.m.*) for the systems 12C+208Pb, 16O+208Pb, 19F+208Pb, 16O+144Sm, and 16O+62Ni. The various results of discussions are summarized in section 1.4.

# Formulation

Nuclear optical model is frequently employed to study collisions of heavy ions. The most important component of scattering and reaction cross-sections is to solve the radial Schro¨dinger equation by using a complex potential. The potential has different parts. The complex nuclear potential *VN*(r), Coulomb potential *VC*(r), and the centrifugal component *Vℓ*(r) are added to provide the necessary effective potential to solve radial equation. While solving the Schro¨dinger equation for this potential, people usually follow the Runge-Kutta method (RK) of numerical integration to get wave function. The derivative of the wave function is also obtained outside the nuclear potential range. This wave function and derivative are further utilized to Coulomb wave function and the derivative of the function to get the results of scattering matrix with an aim to analyse experimental data. However, with this procedure, it is not simple to separate a portion of reaction cross-section from the whole cross-section in order to explain fusion cross-section. As a result, we use a practical but slightly different method to resolve the Schro¨dinger equation [[31–33].](#_bookmark47)

Let's start by carefully examining the s-wave dispersion. A chain of 'n' rectangular potentials, each with an arbitrarily small width 'w', can be thought of as a potential U(r). In reality, a similar method is implied in any numerical integration of a differential equation. After simulating the potential up to the r = Rmax, we have , where is the width of the ith rectangle. Let in the jth region*,* , the strength and width of the potential are denoted by Uj and wj respectively.

The reduced Schro¨dinger equation in this region is given as

|  |  |  |
| --- | --- | --- |
|  |  | (1) |

with the solution given as

|  |  |  |
| --- | --- | --- |
|  |  | (2) |

where the wave number *kj* is defined as for jth segment having width . For both the adjacent segments, we use notation . Here, *m* denotes the particle’s mass and *E* denotes the incident energy. The solution may be expressed in the first three segments near the origin and *r* = 0 as

|  |  |  |
| --- | --- | --- |
|  | , | (3) |
|  |  | (4) |
|  | , | (5) |

Here are coefficients of the wave functions and are arbitrary constants.

When the wave functions is matched with its derivatives near boundary at we have

|  |  |  |
| --- | --- | --- |
|  |  | (6) |
|  |  | (7) |
|  |  | (8) |
| Where |  | (9) |

Similar calculation at the boundary at yields

|  |  |  |
| --- | --- | --- |
|  |  | (10) |
|  |  | (11) |
|  |  | (12) |
| Where |  | (13) |
| and |  | (14) |

This can be generalized for n boundaries to give

|  |  |  |
| --- | --- | --- |
|  |  | (15) |
|  |  | (16) |
|  |  |  |
|  |  | (17) |
| Where |  |  |

Setting the arbitrary constants as and

We get

|  |  |  |
| --- | --- | --- |
|  |  | (18) |
| With |  | (19) |
|  |  |  |
|  |  | (20) |

The function in m-region may be defined in terms of that in the (m-1) area using their cursive nature of the formula for . This allows us to simulate the potential given by *U*(r) by n-step potentials and further create an easy-to-use numerical program for assessment of scattering S-matrix and also the wave function at certain energy. We assume , such that the potential ) is zero for [Eq.1]. Now it is simple to say that the S-wave S-matrix is given by . Total absorption (reaction) cross-section in region is given by

Similarly, is taken as S-matrix of the original potential which is trimmed at . Hence can be taken as the absorption cross section generated in the region . Thus shall give the contribution to the absorption cross section from the region.

We obtain complex conjugate of eq.(1) and rearrange to get

|  |  |  |
| --- | --- | --- |
|  |  | (21) |
|  |  | (22) |
|  |  | (23) |
|  |  | (24) |

We simplify the associated integral using the appropriate wave function and the potential in a specific section and obtain

|  |  |  |
| --- | --- | --- |
|  |  | (25) |
| with |  | (26) |

Considering same width for all segments, i.e .

|  |  |  |
| --- | --- | --- |
|  |  | (27) |

The complex conjugate is represented by an asterisk mark. S-matrix calculation procedure through equations (18), (19) and (20) with a *multistep potential* (MP), an algebraic recursive technique is used for approximation.. This can be programmed easily. The equations (27), (28) and (29) provide a method for study of absorption cross section as the discrete sums of different contributions from different sections.

It is simple to generalize this method to the challenging heavy ion Coulomb nuclear issue for any partial wave. One may use the MP approximation approach for this effective potential to handle higher partial wave problems as the scattering by *VN*(r)+*VC*(r)+ *Vℓ*(r). The complexity of the *rℓ*+1 behaviour of the wave function close to origin in the complex potential scattering is not particularly important for the following reasons. The properly normalized wave function typically attenuates quickly to zero well beyond the origin for absorptive complex potential due to the existence of absorption. As a result, one may begin the S-matrix computation well after r = 0, in which region multistep approximation is more precise. We have confirmed that the cross-section and S-matrix findings produced by our strategy are nearly identical to those produced by traditional methods.

The effective potential has three parts, i.e., in . But the potential is merely Coulombic in the outer region, i.e., in . The last term is centrifugal term for different momentum partial wave .

By using the exact Coulomb wave functions i.e and their derivatives , in the outer region and the wave function and its derivative in the left side of r , we get the expression for partial wave S-matrix as

|  |  |  |
| --- | --- | --- |
|  |  | (28) |
| where |  | (29) |
|  |  | (30) |
|  |  | (31) |

with , which is real at , where the potential is real and

We describe the elastic scattering of a particular system using the preceding formula (30) for. The formula may be used to calculate the  reaction cross-section.

|  |  |  |
| --- | --- | --- |
|  |  | (32) |

As formulated above this is equal to the absorption cross section

|  |  |  |
| --- | --- | --- |
|  |  | (33) |

By taking into account the matching number of segments in the aforementioned summation, it is possible to determine thecontributionofanycomponentwithintherange0*Rmax*totheabsorptionorreactioncrosssection.Withnopotentialdisturbance, this analysis of the collision process's region-by-region absorption results in the wave function that characterises the angular distribution of the elastic scattering data. If someone wants to know how much absorption cross section is there in the region 0 *< r < Rfus*, where*Rfus<Rmax*, the total number of segments to be considered in the summation (35) is . The resulting cross section in given below in Eq.(34). This cross-section refers to fusion in the context of DRM [[14]](#_bookmark29) mentioned earlier.

|  |  |  |
| --- | --- | --- |
|  |  | (34) |

It might be beneficial at this point to briefly explore MP formulation as a numerical approach. In comparison to the trapezoidal rule, Simpson’s rule and spline technique (by using straight-line sections, parabolic sections and cubic polynomials respectively), this method, in a sense, is the most simple approximation to solve differential equations [[34–36].](#_bookmark49)The latter techniques are highly helpful for computing the wave function with more accuracy, but they have the drawback of being difficult to describe algebraically in terms of simple intervals. The computation upto 3-4 significant places after decimal point of wave function and the cross-section, however, is fairly suitable to calculate cross-section given the experimental uncertainties involved. In a study [[37],](#_bookmark50)we compared the numerical results for potentials like Eckart and Ginocchio potentials produced using the MP technique in one dimension with those obtained using the R-K and exact solution, and we discovered that the results agreed upto three significant places. This section has demonstrated how the analytical formulation results in a tidy recursive relation that makes it easier to calculate cross-sections and S-matrices. With this MP formulation, it is possible to transparently estimate the contribution to absorption in various potential segments and examine the wave function nature and the normalization. However, the wave function rises quickly in the calculations for the nucleus-nucleus optical model performed using the R-K method [[38]](#_bookmark51)due to the imaginary potential and should be renormalized at various stages in order to perform calculations of effective phase-shift. This makes it more difficult to estimate the regional contributions to the reactions and the cross-sections [[39].](#_bookmark52)To demonstrate the effectiveness and usefulness of our MP technique, we compute elastic scattering cross-sections and fusion cross-sections in HICs (heavy-ion collisions) and satisfactorily explain the experimental outcomes. The numerical outcomes of the elastic scattering cross sections provided in Ref [38]are validated using conventional optical model techniques.

Table 1: Optical model potential parameters used in the calculations.

*(VB* represents height and *RB* radial position of the Coulomb barrier)

|  |  |  |  |  |  |  |  |  |  |
| --- | --- | --- | --- | --- | --- | --- | --- | --- | --- |
| System *VN* | *rV* | *aV* | *W* | *rW* | *aW* | *rC* | *VB* | *RB* | *Rfus* |
| (MeV) | (fm) | (fm) | (MeV) | (fm) | (fm) | (fm) | (MeV) | (fm) | (fm) |
| 12C+208Pb 125 | 1.31 | 0.320 | 3.0 | 1.325 | 0.25 | 0.90 | 56.7 | 12.16 | 9.9 |
| 16O+208Pb 125 | 1.35 | 0.285 | 2.0 | 1.320 | 0.15 | 1.02 | 73.7 | 12.52 | 8.8 |
| 19F+208Pb 105 | 1.35 | 0.285 | 2.0 | 1.320 | 0.15 | 1.37 | 81.51 | 12.74 | 8.3 |
| 16O+144Sm 100 | 1.295 | 0.365 | 4.0 | 1.250 | 0.15 | 1.1 | 60.25 | 11.46 | 10.0 |
| 16O+62Ni 75 | 1.333 | 0.380 | 6.0 | 1.250 | 0.33 | 1.04 | 30.41 | 10.2 | 7.78 |

# Application

The formulation which is developed in section 1.2 is applied for data analysis of the systems 12C+208Pb, 16O+208Pb, 19F+208Pb, 16O+144Sm and 16O+62Ni. We obtain a reliable explanation of the experimental cross-section magnitudes regarding fusion and elastic scattering. We also find interesting peaks in variation of against the energy Ec.m..

The optical model potential (OMP) in the entrance channel is described by

Here, the analysis involves two nuclei with mass number A1 and A2. Their proton numbers are Z1 and Z2 respectively. We use the form factor in this paper as described below

The symbol VN represents the real part of OMP, whereas, W represents the strength imaginary part. Radius parameters are given as and . The symbols aV and aW are the diffuseness parameters. The Coulomb potential is given by

Where with rC as the Coulomb radius parameter. Thus there are a total of seven parameters, namely, in this potential.

## Elastic scattering cross-section

The angular distribution of elastic scattering can be explained by a number of different sets of potential-descriptive parameters, according to what we know. All seven parameters in our current calculation are independent of energy. While choosing parameter values, we keep the fact in mind that resonance occurs when the imaginary part is kept low. The weak absorption is appreciable when real part is assumed to be deep [[40]](#_bookmark53) and also less diffused [[41]](#_bookmark54) to account for the cross-sections of elastic scattering. The parameter values used for analysis of systems 12C+208Pb, 16O+208Pb, 19F+208Pb, 16O+144Sm , and 16O+62Ni are given in Table 1.1 along with height (VB) and radius (RB) of S-wave barrier in case of each collision system.

1. 12**C+**208**Pb system**

According to the prescription mentioned above, the real portion in case of 12C+208Pb is made more deep with a depth of *VN* = 125 MeV. It is also taken less diffused with a diffuseness parameter of *aV* = 0.32 fm and a radius parameter of *rV* = 1.31 fermi. Again, the corresponding imaginary component also receives the additional parameters *rW* = 1.325 fm, *aW* = 0.25 fm and W = 3.0 MeV. The attractive strength is taken small. The Coulomb radius parameter assumes a value of *rC* = 0.9 fm. Real part of the combined nuclear and Coulomb potentials for an s-wave is shown in **Fig.[1](#_bookmark1)**as a function of radial distance for visual illustration. This demonstrates the repulsive barrier for the 12C+208Pb system with height *VB*= 56.7 MeV and position *RB* = 12.16 fm lowering quickly in the inner side. According to the formulation, ’n’ rectangular potentials, each with a width of 0.008 fm, are used to replicate the potential in the area 0 *< r ≤ Rmax*. The nuclear potential together with its imaginary portion is zero in the area r > Rmax ≈15 fm, leaving the effective potential as the sum of Coulombic and centrifugal term.

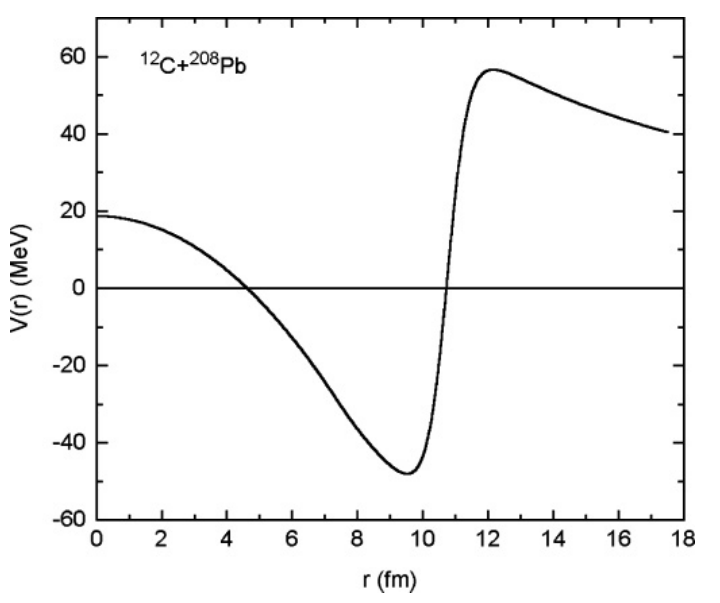


Figure1: Plot of the real part of nuclear plus Coulomb potentials against radial distance with l=0 and with potential parameters*V*0= –125MeV, *rV* = 1.31fm, *aV* = 0.32fm, and *rC* = 0.9 fm for the 12C+208Pb system.

Using the S-matrix given by the expression (1.28), at laboratory energies of 58.9, 60.9, 62.9, 64.9, 74.9, and 84.9 MeV, we get the findings of angular change of differential scattering cross-section. These computed findings are displayed as solid curves in Fig.[2.](#_bookmark2) They are contrasted with the comparable experimental data from Ref.[[1]. The experimental data](#_bookmark16) are depicted in the same figure as solid circles.

It is obvious that the data explanation in each energy case is fairly sound. It should be noted that we utilized exactly the same set of OMP values listed in Table 1.1 to explain the results for energies between 58.9 MeV and 84.9 MeV. In other words, OMP parameter values are independent of energy. We should also point out that the value of *rC*=0.9fm that we are considering is a little lower than the standard value of *rC* = 1.25 fm. This fact is supported by the calculation’s conclusion in Ref. [[42]](#_bookmark55) about Coulomb potentials in HICs, which shows that it has no impact on the outcomes of the elastic collision (cross-section) in our calculation. This lower value of *rC* has been utilized to consistently explain fusion cross-section for the 12C+208Pb system at the low energy. This will be covered in the next section.

1. 16**O+**208**Pbsystem**

We have taken into consideration a deep real potential in this instance as well, with depth *VN* = 125 MeV. The value of diffuseness parameter taken small, i.e., *aV* = 0.285 fm. Other parameters have been given in Table 1.1. The barrier lowering steeply in the inner side with a height of *VB* = 73.7 MeV and a radius of *RB*=12.52 fm for this system can be seen in the graph of real component of nuclear plus Coulomb potentials against radial distance for s-wave in Fig.[3.](#_bookmark3)

Figure [4](#_bookmark4) shows a comparison between our projected differential scattering cross-section values (solid curves) and the corresponding experimental values (shown with solid circles) from Reference [[2]](#_bookmark17)at various laboratory energies, including 80, 83, 88, 90, 96 and 102MeV. It is obvious that the data’s explanation for all energies is sound. Additionally, a single potential is employed for all energies in this case, but with low value of *rC* = 1*.*02fm.

1. 19**F+**208**Pb system**

The potential depth is taken to be *VN* = 105 MeV and the diffuseness parameter is taken to be *aV* = 0.285 fm for the system 19F+208Pb. The depth is taken high, whereas, diffuseness parameter is taken small. Other parameters are mentioned in Table 1.1 with a set of values suitable for analysis. Using the S-matrix given by equation (1.28), the results of angular variation of differential scattering cross-section at energies 88.0, 91.0, 93.0, 96.0, 98.0 and 102.0 MeV in laboratory frame are calculated. In Fig.[5,](#_bookmark5)these computed results are depicted as solid curves, and they are compared with the experimental data from Ref.[[6],](#_bookmark21) which are displayed as solid circles in the same figure. It is obvious that the data’s explanation for all energies is reasonable.

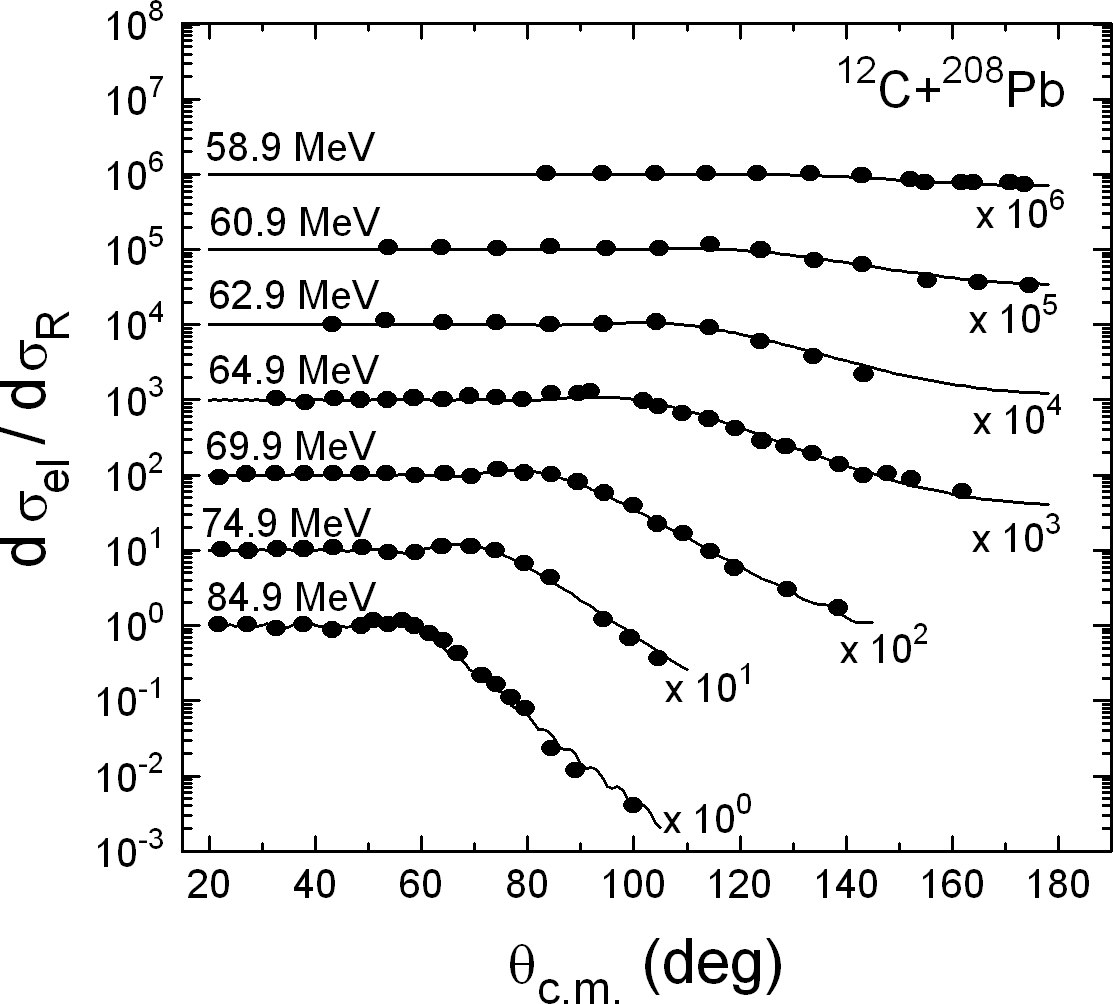


Figure 2: Plotting of scattering cross sections with respect to Rutherford’s cross section for elastic collision system 12C+208Pb. Energies are mentioned in laboratory frame. The angular distribution is given as a function of *θc.m.*. Theoretical calculations are given by solid curves. Solid circles represent experimental data taken from Ref.[[1].](#_bookmark16)

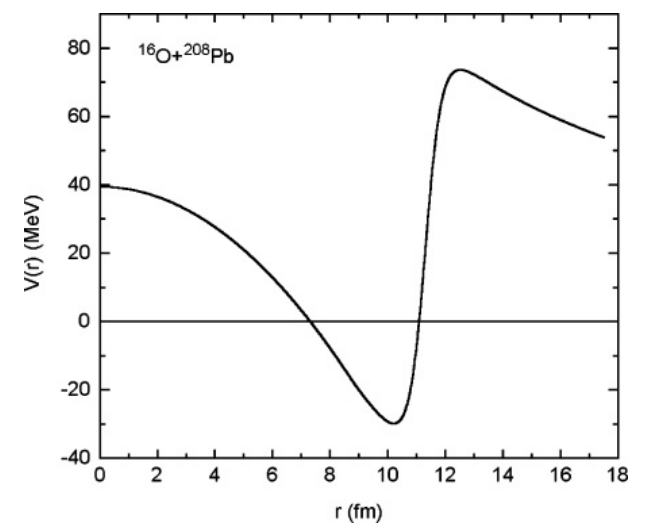


Figure3: Plotting of real part of effective potential (nuclear plus Coulomb potential) against radial distance  for l=0 and with parameters, *V*0= -125MeV, *rV*=1.339fm, *aV*=0.285fm, and *rC*=1.02 fm for 16O+208Pb system.

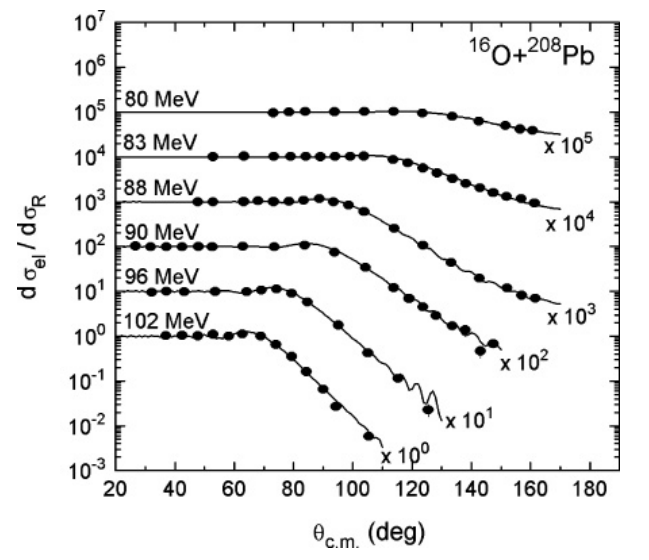


Figure 4: Same as Fig. [2](#_bookmark2) for system 16O+208Pb at 80, 83, 88, 90, 96, and 102 MeV energies in laboratory frame. The solid curves represent theoretical results and solid circles represent experimental data taken from Ref.[[2].](#_bookmark17)

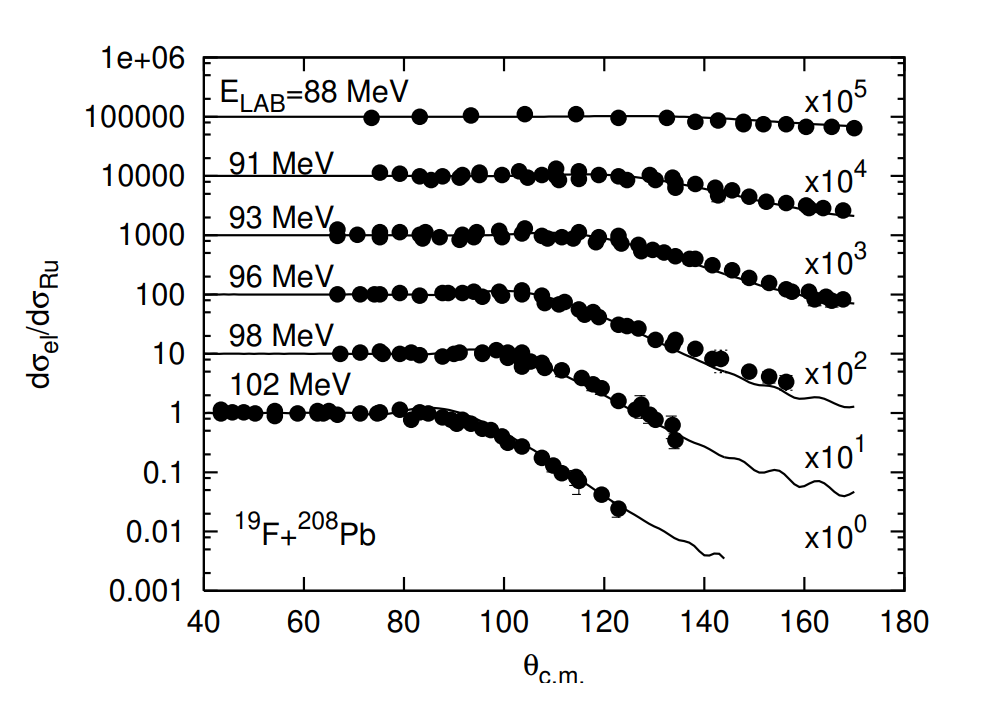
**(v)**16**O+**144**Smsystem**

Similar calculations are done for this system. A deep real potential has been chosen with *VN* = 100.0 MeV and a minimal diffuseness value of *aV* = 0.365 fm in this case as well. Values of other parameter are given in Table 1.1. The plotting in Figure-[6](#_bookmark6) compares the calculated values of differential scattering cross-section (solid curves) with experimental data (solid circles) from Ref.[[8]](#_bookmark23) at various energies, i.e., 66.0, 69.2 and 72.3 MeV. It is seen that the agreement of our theoretical results with experimental data is quite good.

### (vii)16O+62Nisystem

Here, we have used the optical potential parameters of 16O+58Ni system except *rV* = 1.333 fm. In Fig. [7,](#_bookmark7) we compare the differential cross-sections computed theoretically for scattering (given as solid curves) with corresponding experimental data (given as solid circles) obtained from references [[11]](#_bookmark26) at various energies, i.e., 42.0, 48.0 and 54.0 MeV. It is clear from the graphs that the above experimental data are successfully reproduced by our method simultaneously.

Now it is essential to use the same potential in the description of fusion data. The analyses are carried out below for the afore-mentioned systems. The finding of an optical model potential (OMP) which is energy independent is a significant conclusion of our research and discussion.



## 

Figure 5: Angular distribution of elastic scattering cross sections (ratios to Rutherford) as a function of *θc.m.*of 19F+208Pb system at different laboratory energies. The full drawn curves are theoretical results of our optical model calculation. The solid circles are experimental data taken from Ref.[[6].](#_bookmark21)

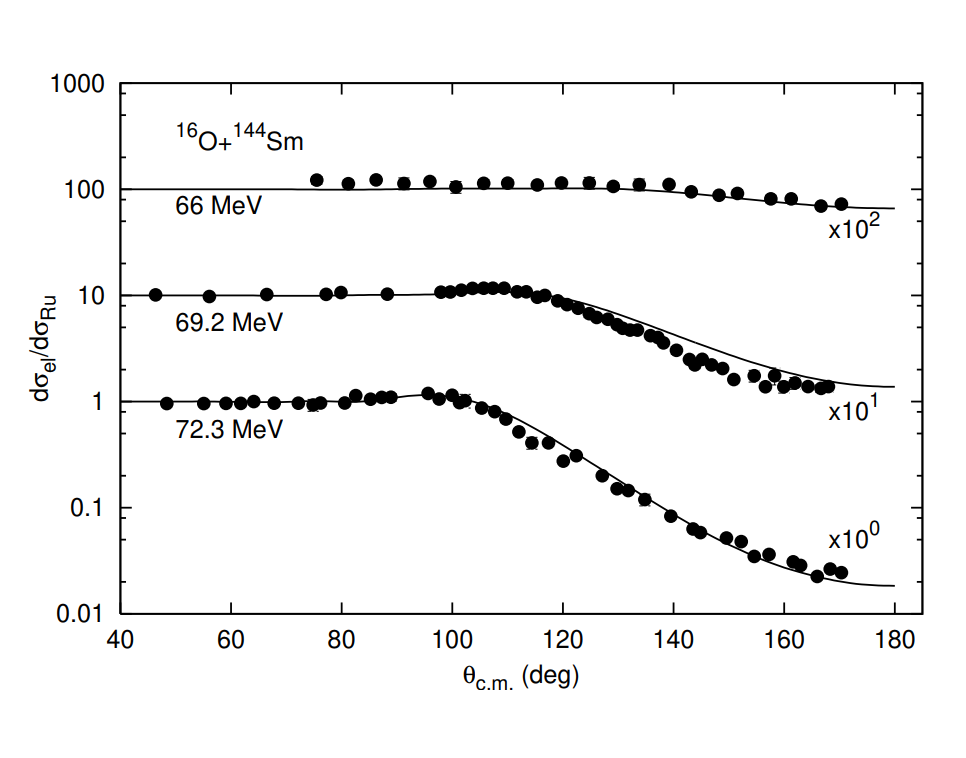


Figure 6: Angular distribution of elastic scattering cross sections (ratios to Rutherford) as a function of *θc.m.* of 16O+144Sm system at differentlaboratory energies. The full drawn curves are theoretical results of our optical model calculation. The solid circles are the experimental data taken from Ref.[[8].](#_bookmark23)

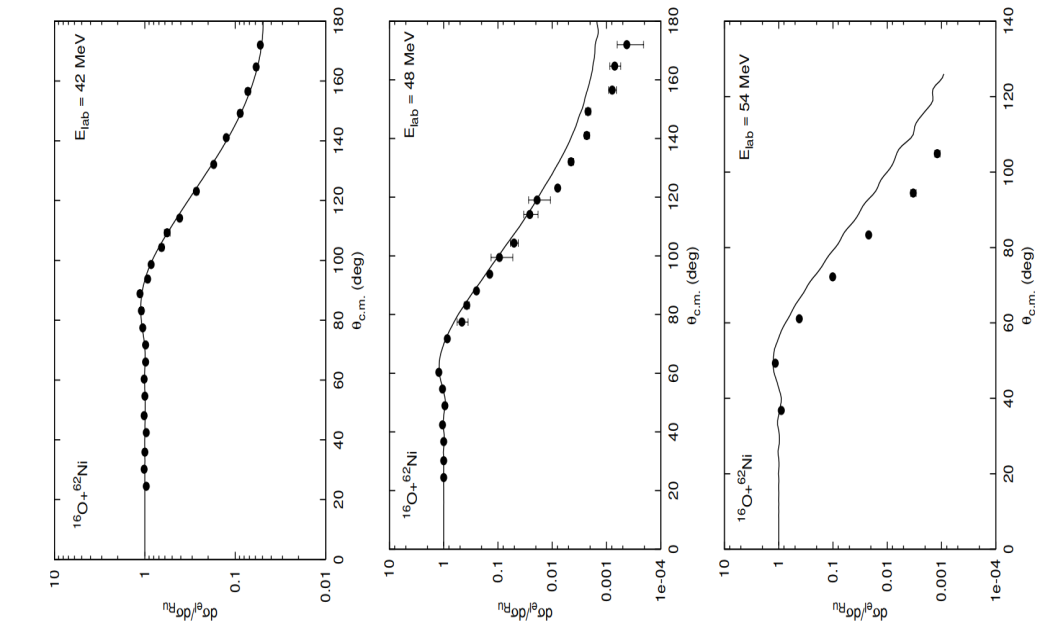


Figure 7: Angular distribution of elastic scattering cross sections (ratios to Rutherford) as a function of *θc.m.*of 16O+62Ni system at differentlaboratory energies. The full drawn curves are theoretical results of our optical model calculation. The solid circles are experimental data taken from Ref.[[11].](#_bookmark26)

## Fusion cross-section

Fusion of the two nuclei is a crucial mechanism. It is actively linked to the elastic collision process in low energy range. It is simple to take into account that cross-section, σfus, of fusion is a portion of the overall reaction cross section σr, when estimating total cross-sections of scattering as well as fusion. However, it is never easy to take a part from σr that precisely accounts for the observed results of fusion at different incidence energy across a large range. In order to compute σfus, we take into account the DRM of Udagawa et al. [[14].](#_bookmark29) The quantity of absorption cross-section inside the inner zone 0 < r < Rfus is what this model refers to as the fusion cross-section. Here Rfus is radial distance which is expected to be smaller than RB. RB is the radial position of S-wave Coulomb barrier in case of a given nuclear system. Ee have calculated in Eq.(1.34) the cross-section values of σfus according to DRM principle.

1. 12**C+**208**Pb system**

We obtain the results of *σfus* for the system 12C+208Pb by using *Rfus*= 9.9 fm. In Fig.8[,](#_bookmark8) the theoretical results (solid curve) so obtained are compared with experimental results (solid circles) collected from Reference [[3].](#_bookmark18) It is evident from the graphs that the data fairly match throughout the whole energy range, i.e., from E*c.m.*= 50 MeV to 75 MeV. The OMP parameters, whose values explain scattering data in Fig.[2,](#_bookmark2) have not been altered in order to achieve this fitting. The values of *Rfus*= 9.9 fm utilized in our computation are smaller than Coulomb radius *RB*=12.16 fm, which is required by the permissible state of reality. The present successful description of scattering and fusion cross-sections is more important because of the fact [[3,](#_bookmark18)[26]](#_bookmark41) that Woods-Saxon potential fails to explain.

1. 16**O+**208**Pbsystem**

Remarkable accomplishment is attained in case of 16O+208Pb system, where our calculated results (solid curve) are able to reproduce experimental results as indicated by solid dots of *σfus* taken from Reference [[4]](#_bookmark19) for the given range of energy from *Ec.m.*=68 MeV to *Ec.m.*= 86 MeV. Here, fusion radius, *Rfus*=8.8 fm is small compared to Coulomb radius *RB* = 12.52 fermi. Unlike in case of 12C+208Pb, inthis16O+208Pbsystem, we need slight modification of nuclear radius parameter *rV*=1.35 fm (Table1.1). This is taken for analyzing scattering data and take *rV* = 1.339 fm for the fitting of measured *σfus* data.

1. 19**F+**208**Pbsystem**

In this case, we have used the fusion radius *Rfus*=8.3 fm. This radius is smaller than the Coulomb radius, which is taken as *RB*=12.74 fm. We compare, in Fig.10,our calculated cross-sections of fusion with the corresponding experimental cross-sections obtained from Ref. [[7].](#_bookmark22) In order to get a good fitting in fusion cross section, we modify the nuclear radius parameter slightly from *rV* = 1.35 fm to *rV* = 1.356 fm for successful analysis of experimental elastic scattering results.

**(iv)**16**O+**144**Smsystem**

When our calculated results (solid curve) are compared with the experimental results shown by solid dots of cross-section *σfus* taken from Ref. [[9]](#_bookmark24) for the 16O+144Sm system as shown in Fig.[11,](#_bookmark10)a similar result is achieved. The fusion radius in this instance, *Rfus*=10.0 fm, is less than the Coulomb radius, *RB*=11.46 fm. The nuclear radius parameter, *rV*, which was previously set at 1.145 fm (Table 1.1) for the analysis of scattering data must now be changed to 1.295 fm for fair reproduction of observed *σfus* data.

### (v) 16O+62Ni collision system

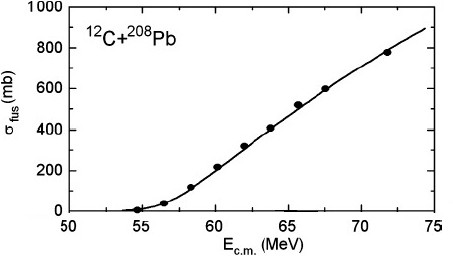
We have taken *Rfus*=7.78 fm which is less than *RB*=10.20 fm. The calculated values (as shown by solid curve) of fusion cross-section is compared with the experimental results (given as solid circles) taken from Ref.[[11]](#_bookmark26)in Fig.[12.](#_bookmark11)It is clear from the figure that the theoretical results is in excellent agreement with the experimental data. To achieve this we have slightly changed the *rV* value from 1.333fm (used for elastic scattering calculation) to 1.313fm.

Figure 8:Variationof fusion cross section *σfus*as a functionof energy E*c.m.*forthe 12C+208Pb system. The solid and continuous curve represents our theoretical calculation. Experimental data which are shown by solid circles are found from Ref.[[3].](#_bookmark18)

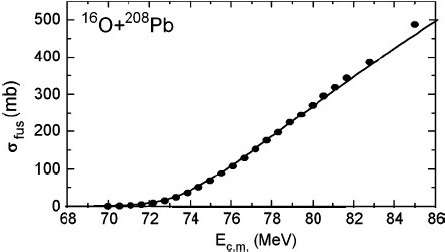


Figure9:Variation of cross-section *σfus* as a function of E*c.m.* for the 16O+208Pb system. The solid curve represents our theoretical calculation. Experimental values which are shown by dark circles are found from Ref.[[4].](#_bookmark19)

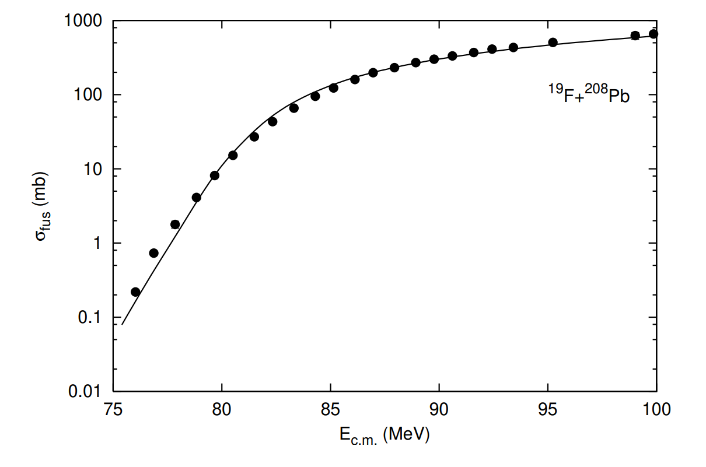


Figure10: Variation of fusion cross-section *σfus* as a function of energy E*c.m.* for the 19F+208Pb system. The continuous curve represents our theoretical calculation. Experimental data are represented by dark circles. The experimental data are found from Reference [[7].](#_bookmark22)

The fusion radius *Rfus* values we utilized in our calculations are, however, smaller in all of these systems than the corresponding Coulomb radius *RB* values. This indicates unequivocally that fusion is an internal phenomena whereas scattering and other distant, less absorptive direct reaction mechanisms are responsible for the surface occurrence. We should note that for a given system, there may be many sets of possible Woods-Saxon parameters that describe elastic scattering data in a manner that is comparable. Since elastic scattering is a surface phenomenon, all sets of potential parameter sets produce Coulomb barriers with the same height *VB* and fixed radial position *RB*, but different sets produce different depths and slopes of effective potential of interior portion *r < RB*. However, the fusion is an internal event that is defined by absorption in this area, and the values of the radius parameter *Rfus*, which is located in the region 0*<r<RB*, are responsible for the corresponding cross-section. The value of *Rfus*will be determined by a set of suitable parameters employed in the study. Therefore, it may have different values for different sets of potential. We do not need to alter the value of *Rfus* as a function of energy for study of *σfus* at various incidence energies as we have chosen a single potential for the description of both the fusion and elastic cross-sections. *Rfus*’s energy-independent character is essential since it makes the product Ec.m σfus in the result of d2(Ec.m σfus)/dE2 , which is described in the next section.

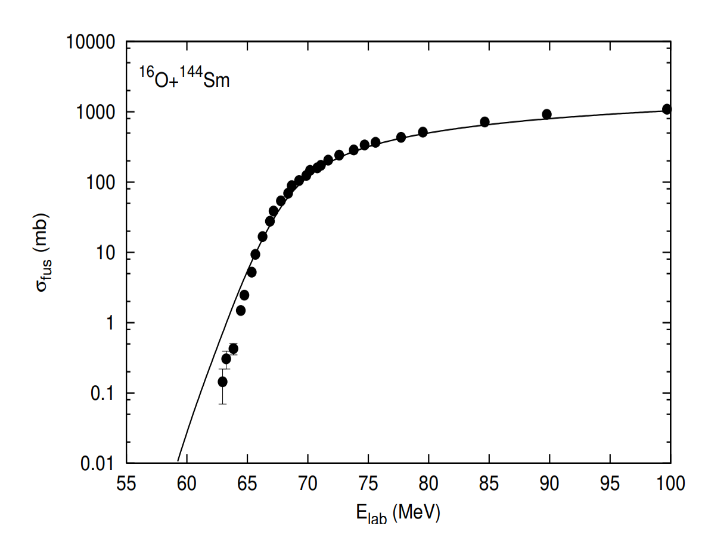


Figure11:Variationoffusioncrosssection*σfus*asafunctionofenergyE*lab*forthe16O+144Smsystem.Thefullcurverepresentsourcalculatedresults.Theexperimentaldatashownbysolidcirclesareobtainedfrom[[9].](#_bookmark24)

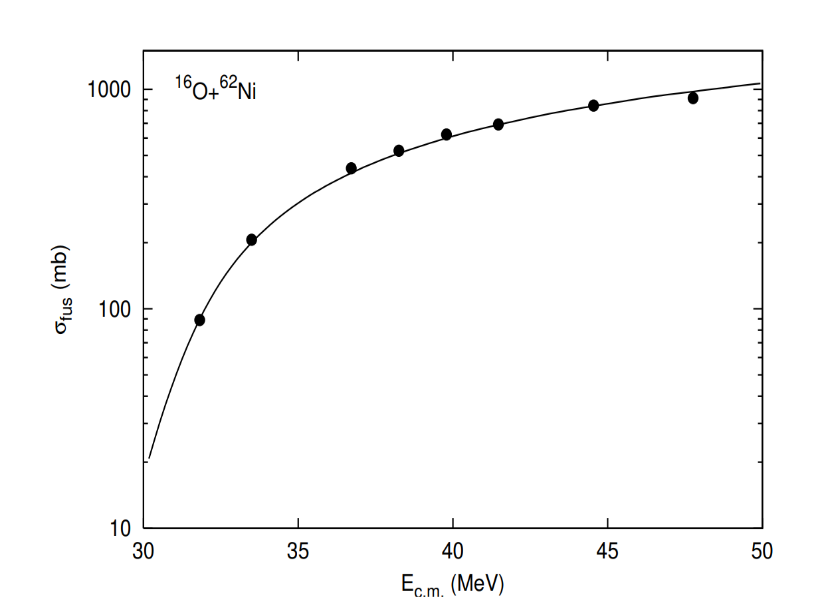
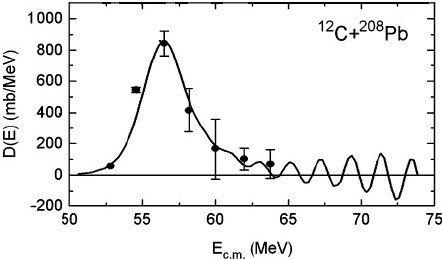


Figure 12: Variation of fusion cross section *σfus*as a function of energy E*c.m.*for the 16O+62Ni system.The full curve represents our calculatedresult. Theexperimentaldatashownbysolidcirclesareobtainedfrom[[11].](#_bookmark26)

Figure13:Variation of against energy Ec.m for the results of for 12C+208Pb system. The full curves represent our calculated results. The experimental data shown by solid dots are obtained from [3].

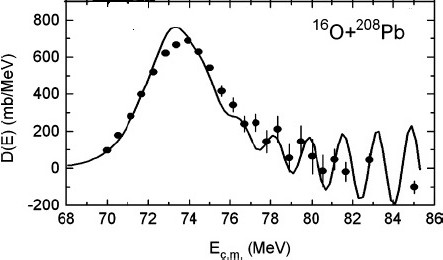


Figure14: Variation of against energy Ec.m for the results of for 16O+208Pbsystem. The full curves represent our calculated results. The experimental data shown by solid dots are obtained from [4].

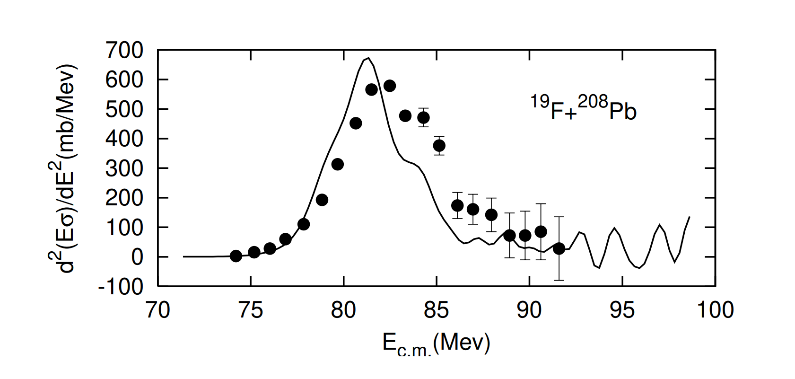
## Explanation of D(E*c.m.*)

Fig.8 to Fig.12 show the values of *σfus* which are obtained from experiment and theory against energy. They do not exhibit any kind of structure. Therefore, nothing further can be inferred about the potential physical incidents that could be contributing to the fusion process from this logical interpretation of the monotonically evolving data. The identical result of *σfus* is reported in a different way as follows in order to provide some insight into these processes. One can extract values of a quantity that is the second derivative of the product Ec.m σfus denoted by with respect to Ec.m. Thus, the point difference formula is given by

|  |  |  |
| --- | --- | --- |
|  |  | (35) |

Where σ – , σ and σ+ indicate fusion cross sections σfus at centre of mass energies E – ΔE, E and E + ΔE, respectively with energy step size ΔE. The function D (Ec.m) is referred to as barrier distribution function [5, 9, 43].We get the amount *D*(*Ec.m.*) against energy *Ec.m*. from our computed results of *σfus* using the formula (1.35). For the12C+208Pb system, we present our results in Fig. [13](#_bookmark12) as a solid curve and then compare them with the relevant experimental data(solid circles). It can be observed that our calculation accurately reproduces the major peak as well as a few other minor peaks in the higher energy range.

Similarly, in Figs. 14, 16, and 17 we remarkably explain high oscillatory structures of the experimental data of *D*(*Ec.m.*) in 16O+208Pb, 19F+208Pb, 16O+92Zr, 16O+144Sm, 16O+62Ni, and 6Li+209Bi systems.

Figure15: Variation of against energy Ec.m for the results of for 19F+208Pb system. The full curves represent our calculations. The experimental data shown by solid dots are by solid circles are obtained from [7].

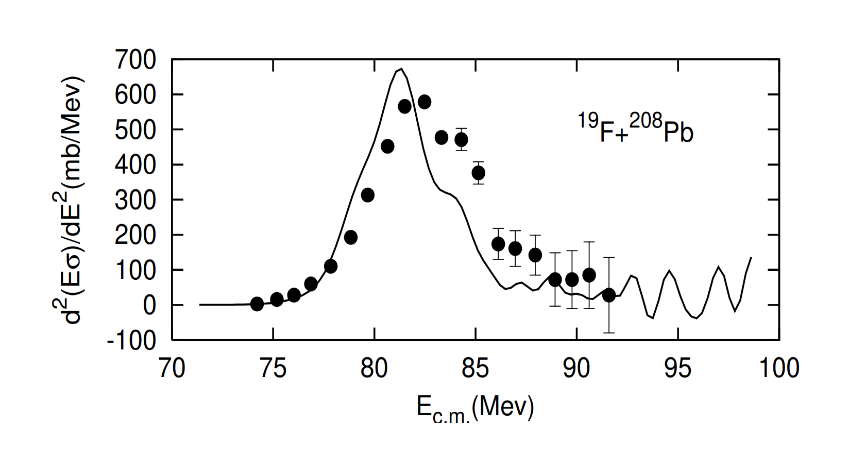


Figure 16 : Variation of *D*(*E*) as a function of energy *Elab* corresponding to results of *σfus* for 16O+144Sm system. The full curve represents our calculated result. The experimental data (solid circles) are obtained from [[10].](#_bookmark25)

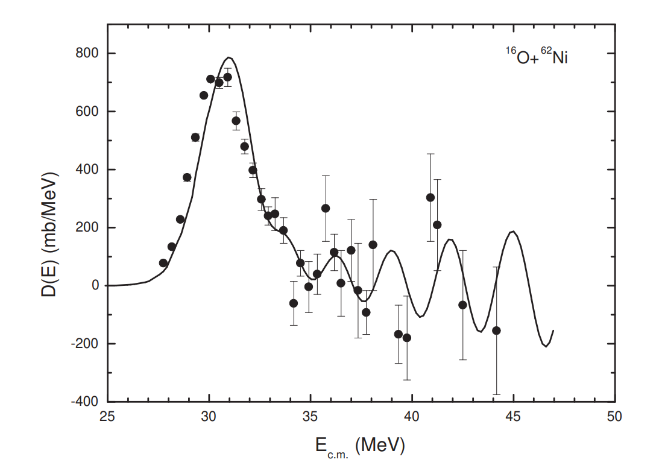


Figure17:Variation of against energy Ec.m for results of for 16O+62Nisystem. The full curves represent our calculated results. The experimental data shown by solid dots are by solid circles are obtained from [7].

More crucially, the unfavourable character of some dips in the zone of higher-energy is fairly explained. The discovery [[4]](#_bookmark19) that the more microscopic CC calculation [[44]](#_bookmark57) for fusion has failed to explain the results of *D*(*Ec.m.*) in the 16O+208Pb system [[4,](#_bookmark19)[23]](#_bookmark38) raises the significance of this successful explanation. It’s worth noting that there are three ways to disrupt the oscillatory structures in all of the systems mentioned above. The ways are as follows :

(i) by raising the value of the imaginary component W’s strength,

(ii) by taking into account a higher value of *rC*, and

(iii) by widening step size ∆*E* for the differentiation using formula (1.35).

In order to provide an appropriate explanation for the elastic scattering data as well as an explanation for the observed results of cs *σfus* and *D*(*Ec.m.*), when the values of W and *rC* are fixed.

# Summary and conclusion

Analytical solutions to the Schro¨dinger equation for the interaction of two nuclei with an optical potential of composite fashion result in a recursive mathematical formula for the scattering matrix. The formula (analytical) for the absorption cross-section has been established to take into consideration the reaction cross-section by using the same potential and wavefunction. The formulation is used for the colliding systems, namely, 12C+208Pb, 16O+208Pb, 19F+208Pb, 16O+144Sm , and 16O+62Ni to analyze of the following experimental results consistently.

1. Angular fluctuations of differential cross-sections of elastic scattering at various incident energies near the Coulomb barrier.
2. Behaviour of cross-sections ( *σfus* ) of fusion against energy throughout a large range spanning Coulomb barrier area.
3. Outcome of the distribution quantity .

The key characteristics that result from this analysis may be summed up as follows.

(a) Experimental data regarding elastic-scattering at various energies of projectile could be well explained by a single potential based on the format of Woods-Saxon potential without dependence on energy. The real part the complex potential considered should have larger depth, but smaller diffuseness, whereas, the imaginary part assumes weak strength for less absorption.

(b) A key component of the computation is the estimation of the reaction cross-sectional area that will be used for estimating fusion cross-sections by using the stepwise absorption approach. This procedure of dividing up the total cross-section is natural because neither the extraction process nor the division of the imaginary part ever requires additional energy.

(c) The results of cross-section are presented in a form described as, by the technique of point difference formula. The new form shows peculiar peak structure when plotted against energy . Our calculated outcomes for *σfus* presented in the aforementioned manner describe this outcome with highs and lows with a surprising degree of success.

(d) It has been found that resonance states can develop when two nuclei collide because of theoptical potential’s weakly absorptive property indicated in item (a) above. Following that, it comes out that these resonances regulate the interesting oscillatory phenomena found in *D*(*Ec.m.*) mentioned in point (c).

# References

* 1. S.Santra,P.Singh,S.Kailas,A.Chatterjee,A.Shrivastava,andK.Mahata,*Phys.Rev.***C64**,024602(2001).
  2. F.Videbaek,R.B.Goldstein,L.Grodzins,S.G.Steadman,T.A.Belote,andJ.D.Garrett,*Phys.Rev.***15**,954(1977).
  3. A.Mukherjee,D.J.Hinde,M.Dasgupta,K.Hagino,J.O.Newton,andR.D.Butt,*Phys.Rev.***C75**,044608(2007).
  4. C.R.Morton,A.C.Berriman,M.Dasgupta,D.J.Hinde,J.O.Newton,K.Hagino,andI.J.Thompson,*Phys.Rev.***C60**,044608(1999).
  5. M.Dasgupta,D.J.Hinde,N.Rowley,andA.M.Stefanini,*Annu.Rev.Part.Sci.***48**,401(1998),andreferencestherein.
  6. C.J.Lin,J.C.Xu,H.Q.Zhang,Z.H.Liu,F.Yang,andL.X.Lu,*Phys.Rev.***C63**064606(2001).
  7. D.J.Hinde,A.C.Berriman,M.Dasgupta,J.R.Leigh,J.C.Mein,C.R.Morton,andJ.O.Newton,*Phys.Rev.***C60**054602(1999).
  8. D. Abriola, D. DiGregorio, J. E. Testoni, A. Etchegoten, J. O. Fernandez, A. M. J. Ferrero, S. Gil, A. O. Macchiavelli, A. J.Pacheco,andJ.Kittl,*Phys.Rev.***C39**546(1989).
  9. J.R.Leigh,M.Dasgupta,D.J.Hinde,J.C.Mein,C.R.Morton,R.C.Lemmon,J.P.Lestone,J.O.Newton,H.Timmers,J.

X.Wei,andN.Rowley,*Phys.Rev.***C52**,3151(1995).

* 1. K.Hagino,N.Takigawa,andS.Kuyucak,*Phys.Rev.Lett.***79**2943(1997).
  2. N.Keeley,J.A.Christly,N.M.Clarke,B.R.Fulton,J.S.Lilley,M.A.Nagarajan,I.J.Thompson,*Nucl.Phys.***A582**314(1995).
  3. N.Keeley,J.S.Lilley,J.X.Wei,M.Dasgupta,D.J.Hinde,J.R.Leigh,J.C.Mein,C.R.Morton,H.Timmers,andN.Rowley,*Nucl.Phys.***A628**1(1998).
  4. Y.EisenandZ.Vager,*Nucl.Phys.***A187**,219(1972).
  5. T.Udagawa,B.T.Kim,andT.Tamura,*Phys.Rev.***C32**,124(1985).
  6. T.Udagawa,S.-W.Hong,andT.Tamura,*Phys.Rev.***C32**,1435(1985).
  7. J.R.BirkelundandJ.R.Huizenga,*Annu.Rev.Nucl.Part.Sci.***33**,265(1983).
  8. S.G.SteadmanandM.J.Rhoades-Brown,*Anuu.Rev.Nucl.Part.Sci.***36**,649(1986).
  9. G.R.Satchler,M.A.Nagarajan,J.S.Lilley,andI.J.Thompson,*Phys.Rev.***C41**,1869(1990).
  10. M.A.NagarajanandG.R.Satchler,*Phys.Lett.***B173**,29(1986).
  11. C.H.DassoandS.Landowne,*Comput.Phys.Commun.***46**(1987)187;M.Dasgupta,A.Navin,Y.K.Agrawal,C.V.K.Baba,H.C.Jain,M.L.Thingan,andA.Roy,*Nucl.Phys.***A539**,351(1992).
  12. I.J.Thompson,*Comput.Phys.Rep.***7**,167(1998).
  13. K.Hagino,N.Rowley,andA.T.Kruppa,*Comput.Phys.Commun.***123**,143(1999).
  14. HenningEsbensenandSerbanMisicu,*Phys.Rev.***C76**,054609(2007).
  15. M.Dasgupta,D.J.Hinde,A.Diaz-Torres,B.Bouriquet,CatherineI.Low,G.J.Milburn,andJ.O.Newton,P*ˆ*hys.Rev.Lett.

**99**,192701(2007).

* 1. R.N.Sagaidak,S.P.Tretyakova,S.V.Khlebnikov,A.A.Ogloblin,N.Rowley,andW.H.Trzaska,*Phys.Rev.***C76**,034605(2007).
  2. J.O.Newton,R.D.Butt,M.Dasgupta,D.J.Hinde,I.I.Gontchar,C.R.Morton,andK.Hagino,*Phys.Rev.***C70**,024605(2004).
  3. D.J.Hinde, M.Dasgupta, C.R.Morton, A.C. Berriman,R.D. Butt,andJ.O.Newton, *Nucl.Phys.***A654**,864c (1999);D.

J.Hinde,C.R.Morton,M.Dasgupta,J.R.Leigh,J.C.Mein,andH.Timmers,*Nucl.Phys.***A592**,271(1995).

* 1. M.R.Spinella,J.E.Testoni,O.Dragun,andH.D.Matra,*Nucl.Phys.***A687**,385(2001).
  2. G.R.Satchler,M.A.Nagarajan,J.S.Lilley,andI.J.Thompson,*Ann.Phys.***178**,110(1987).
  3. M.L.Halbert,J.R.Beene,D.C.Hensley,K.Honkanen,T.M.Semkow,V.Abenante,D.G.Sarantites,andZ.Li,*Phys.Rev.*

**C40**,2558(1989).

* 1. BasudebSahu,G.S.Mallick,B.B.Sahu,S.K.AgarwallaandC.S.Shastry,*Phys.Rev.***C77**,024604(2008).
  2. R.R.Swain,C.Dash,P.Mohanty,B.B.Sahu*Int.J.Mod.***E29**,2050016(2020).
  3. C.Dash,R.R.Swain,G.Tripathy,I.Naik,B.B.Sahu.*ChinesePhysics***C46**,124104-1-8,(2022).
  4. R.H.LandauandM.J.P.Mejia,*ComputationalPhysics*(JohnWileyandSons,Newyork,1997).
  5. F.R.LoscalzoandT.D.Talbot,*Bull.Amer.Math.Soc.***73**,438(1967).

[36]G.Micula,*Math.Comp.***27**,807(1973).

1. BasudebSahu,BidhubhusanSahu,andSantoshK.Agarwalla,*Pramana:J.Phys.***70**(2007)27.
2. M.A.Melkanoff,T.Sawada,andJ.Raynal,*MethodsinComputationalPhysics,Vol.VI,NuclearOpticalModelCalculations*(AcademicPress,NewYork,1966).
3. C.S.ShastryandY.K.Gambhir,*Phys.Rev.***C28**,1109(1983).
4. M.P.Nicoli,F.Haas,R.M.Freeman,N.Aissaoui,C.Beck,A.Elanique,R.Nouicer,A.Morsad,S.Szilner,Z.Basrak,M.E.Brandan,andG.R.Satchler,*Phys.Rev.***C60**,064608(1999).
5. S.C.Pieper,M.J.Rhoades-Brown,andS.Landowne,*Phys.Lett.***B162**,43(1985).
6. R.M.DevriesandM.R.Clover,*Nucl.Phys.***A243**,528(1975).
7. A.B.BalantekinandN.Takigawa,*Rev.Mod.Phys.***70**,77(1998).
8. N.Rowley,G.R.Satchler,andP.H.Stelson,*Phys.Lett.***B254**,25(1991).
9. B.Sahu,L.Satpathy,andC.S.Shastry,*Phys.Lett.***A303**,105(2002).
10. B.Sahu,S.K.Agarwalla,andC.S.Shastry,*Nucl.Phys.***A713**,45(2003).